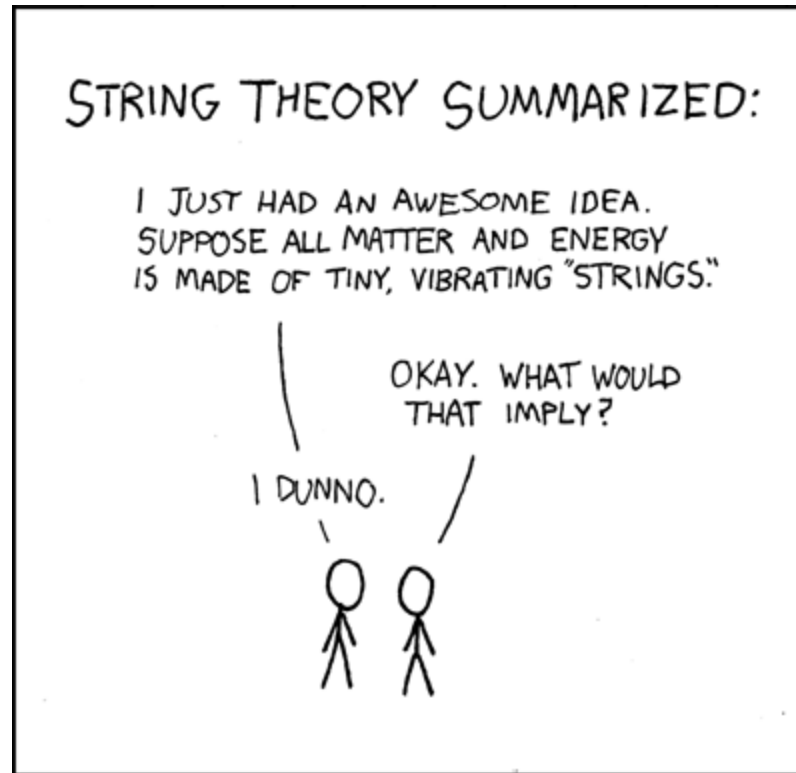


String Theory:

A Model Beyond Popular Physics



Kenny Wunder

Mississippi State University

TAMU Cyclotron

Types of string theory

- Bosonic string theory
- Type I string theory
- Type II-A string theory
- Type II-B string theory
- Heterotic string theory



M-theory

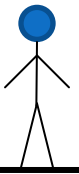
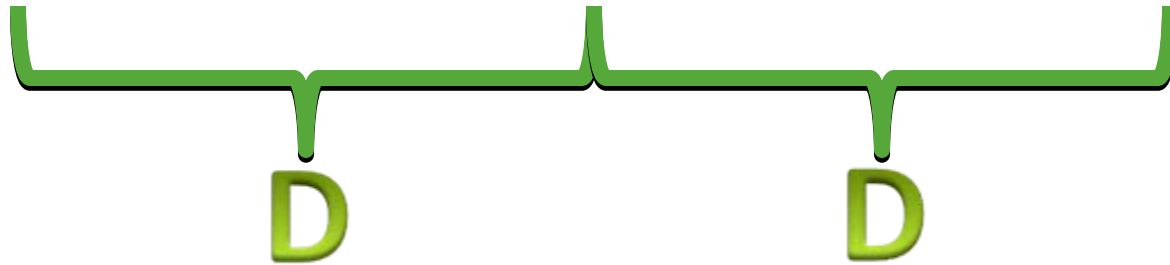
The **big** ideas *(in **bosonic string theory**)*

- On geometry...
 - Identification
 - Compactification
- D-branes
- Evolution of the action
- Parameterizing made easy
 - The light cone gauge*
- Becoming particles

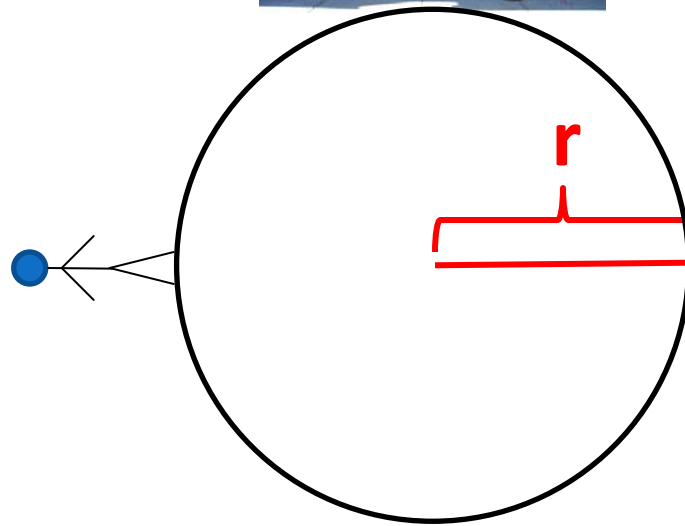
Some things about geometry: Identification



Some things about geometry: Identification

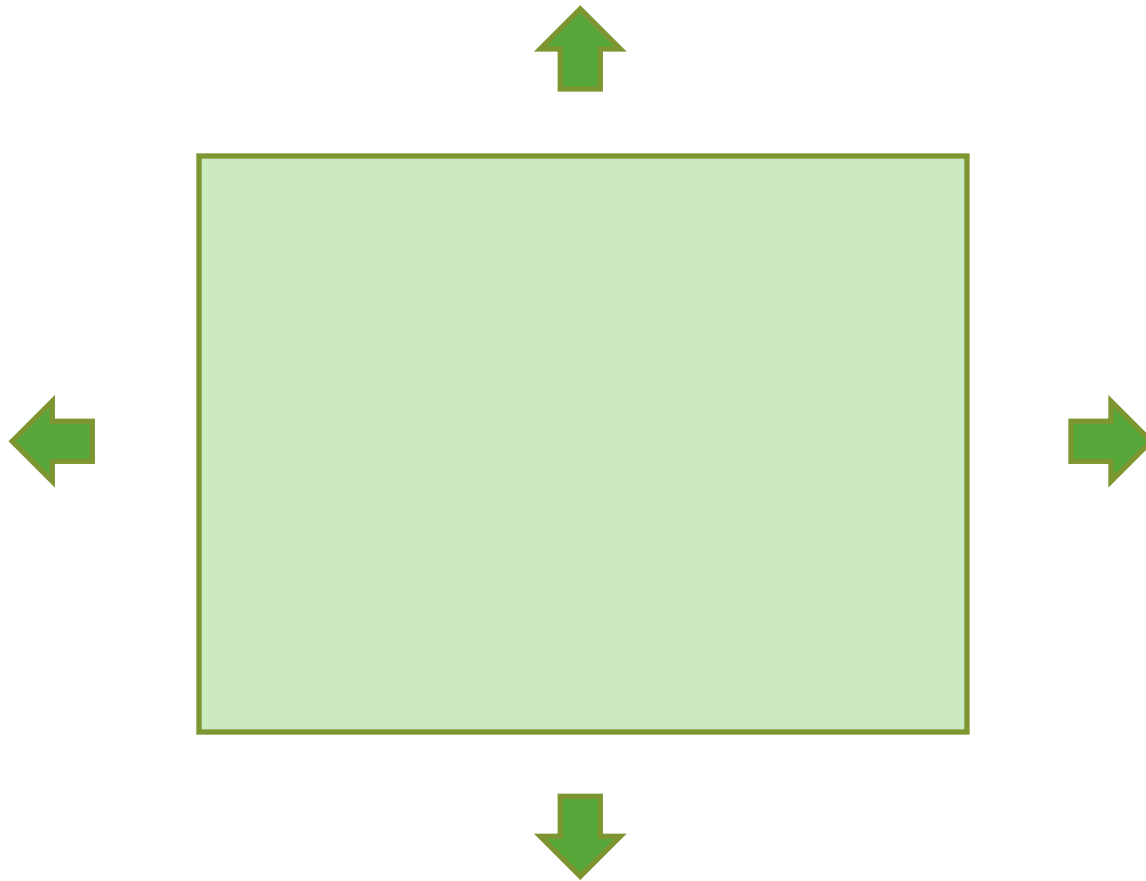


Some things about geometry: Identification

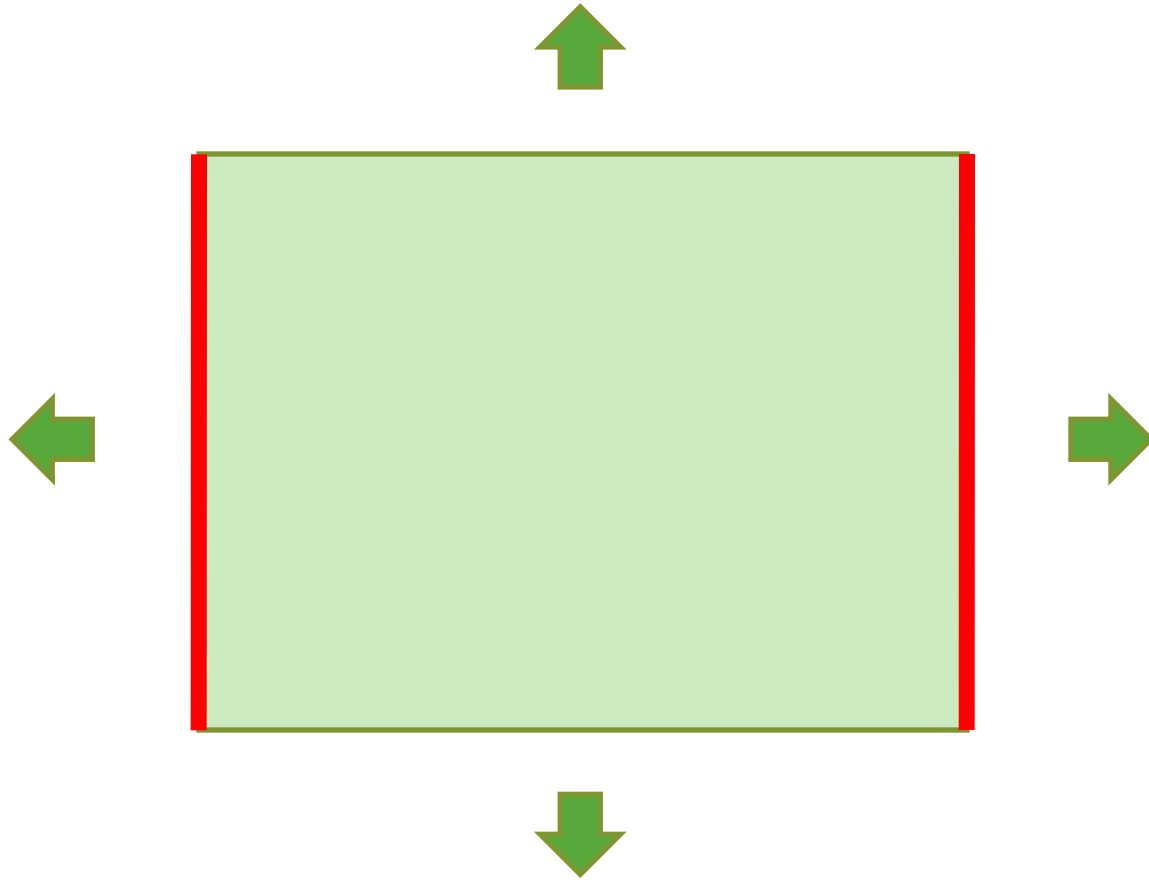


$$D = 2\pi r$$

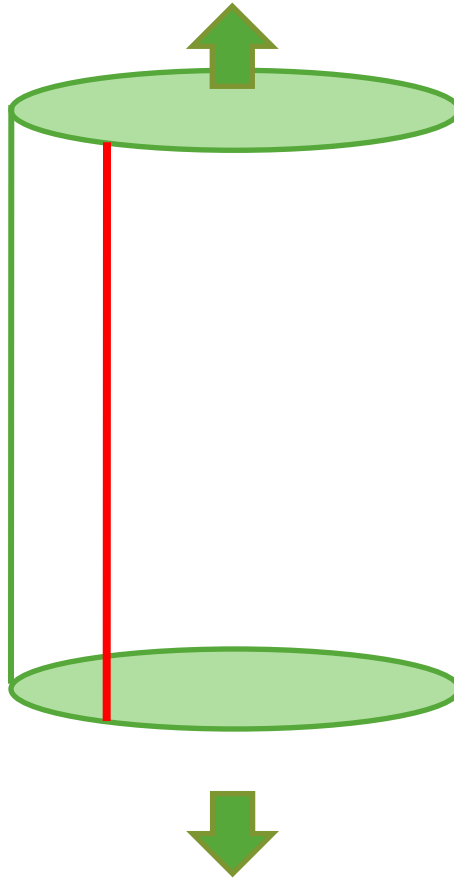
Some things about geometry: Compactification



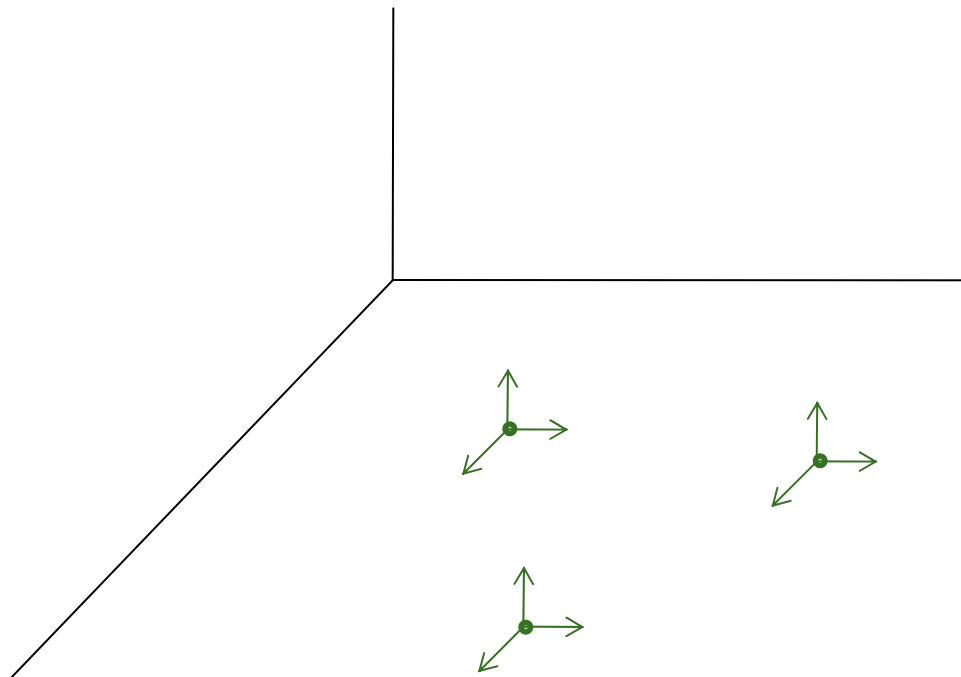
Some things about geometry: Compactification



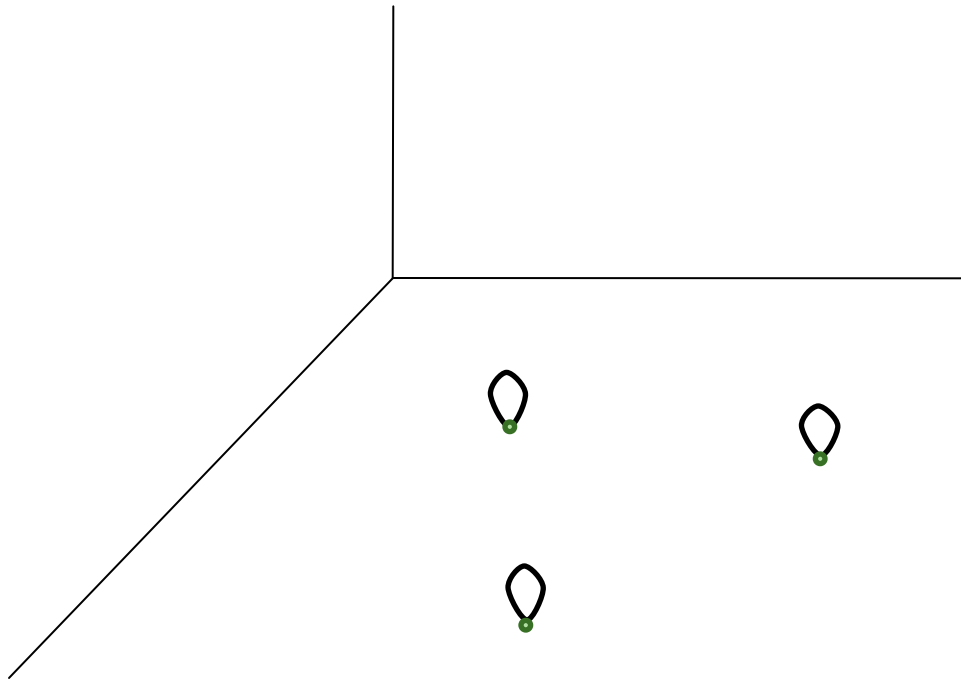
Some things about geometry: Compactification



Some things about geometry: Compactification



Some things about geometry: Compactification



Some things about geometry: Compactification



Boundary conditions and D-branes

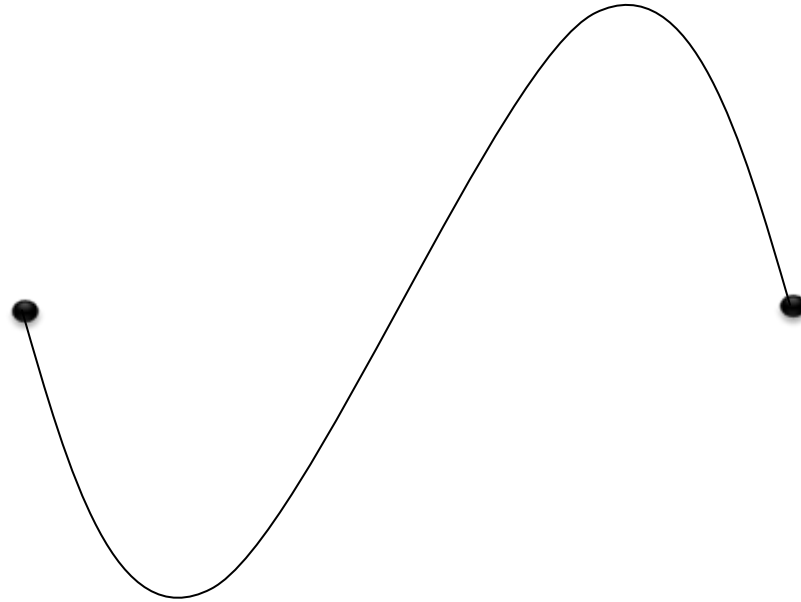
- Neumann boundary conditions

$$\frac{\partial y}{\partial x}(t, 0) = 0 = \frac{\partial y}{\partial x}(t, a)$$

- Dirichlet boundary conditions

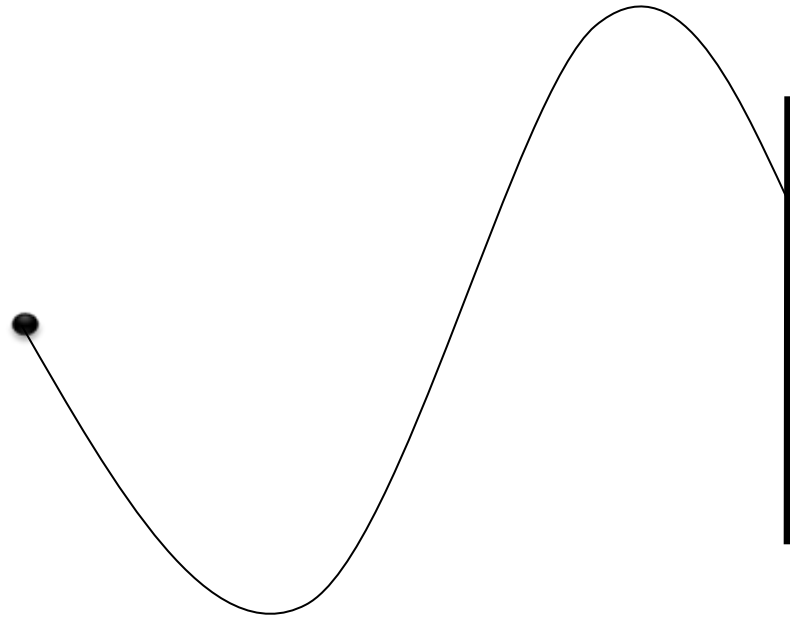
$$y(t, 0) = 0 = y(t, a)$$

Boundary conditions and D-branes



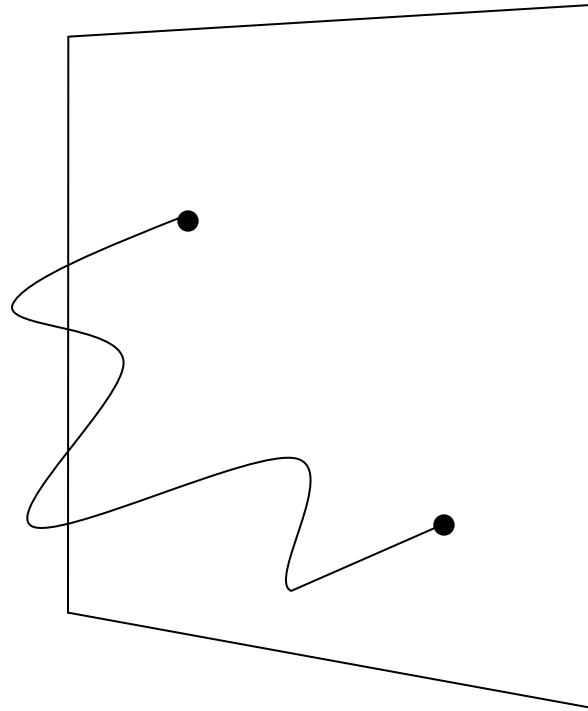
D0-brane

Boundary conditions and D-branes



D1-brane

Boundary conditions and D-branes



D2-brane

The classical action

The classical point particle

Energy $\frac{1}{2}m\dot{x}^2 - V(x)$ Considerations $\frac{d}{dt}$

+ Principle $\int_{t_i}^{t_f} \left\{ \frac{1}{2}m\dot{x}^2 - V(x) \right\} dt$ of Least Action

$$\delta S = \int_{t_i}^{t_f} \delta x \left(-m\ddot{x} - V'(x) \right) dt = 0 \rightarrow m\mathbf{F} = m\mathbf{\ddot{a}}(x)$$

The classical action

The classical string

$$L(t) = \int_0^a \left[\frac{1}{2} (\mu_0) \left(\frac{\partial y}{\partial t} \right)^2 - \frac{1}{2} T_0 \left(\frac{\partial y}{\partial x} \right)^2 \right] dx$$

$$S = \int_{t_i}^{t_f} L(t) dt = \int_{t_i}^{t_f} dt \int_0^a dx \left[\frac{1}{2} (\mu_0) \left(\frac{\partial y}{\partial t} \right)^2 - \frac{1}{2} T_0 \left(\frac{\partial y}{\partial x} \right)^2 \right]$$

$$\delta S = \int_{t_i}^{t_f} dt \int_0^a dx \left[\frac{\partial}{\partial t} \left(\mu_0 \frac{\partial y}{\partial t} \delta y \right) - \mu_0 \frac{\partial^2 y}{\partial t^2} \delta y + \frac{\partial}{\partial x} \left(-T_0 \frac{\partial y}{\partial x} \delta y \right) + T_0 \frac{\partial^2 y}{\partial x^2} \delta y \right]$$

The classical action

The classical string

$$\mathcal{L}\left(\frac{\partial y}{\partial t}, \frac{\partial y}{\partial x}\right) = \frac{1}{2}(\mu_0)\left(\frac{\partial y}{\partial t}\right)^2 - \frac{1}{2}T_0\left(\frac{\partial y}{\partial x}\right)^2$$

$$\mathcal{P}^t \equiv \frac{\partial \mathcal{L}}{\partial \dot{y}} = \mu_0 \frac{\partial y}{\partial t} \qquad \mathcal{P}^x \equiv \frac{\partial \mathcal{L}}{\partial y'} = -T_0 \frac{\partial y}{\partial x}$$

The classical action

The classical string

$$\delta S = \int_{t_i}^{t_f} dt \int_0^a dx \left[\frac{\partial}{\partial t} \left(\mu_0 \frac{\partial y}{\partial t} \delta y \right) - \mu_0 \frac{\partial^2 y}{\partial t^2} \delta y + \frac{\partial}{\partial x} \left(-T_0 \frac{\partial y}{\partial x} \delta y \right) + T_0 \frac{\partial^2 y}{\partial x^2} \delta y \right]$$



$$\delta S = \int_{t_i}^{t_f} dt \int_0^a dx \left[\mathcal{P}^t \delta \dot{y} + \mathcal{P}^x \delta y' \right]$$

The classical action
The classical string

Again:
 $\frac{\partial \mathcal{P}}{\partial t}$ $\frac{\partial \mathcal{P}}{\partial x}$
Energy Considerations
+ Principle of Least Action

the **wave equation!!**

The relativistic action

The relativistic point particle

- Using Space-Time Distance $\sqrt{c^2 dt^2 - dx^2}$

+ Energy Considerations $\int_{t_i}^{t_f} dt \sqrt{c^2 - v^2}$

So, Relativistic Action $\int_{t_i}^{t_f} dt \sqrt{c^2 - v^2}$

The relativistic action

The relativistic point particle

Beginning with the

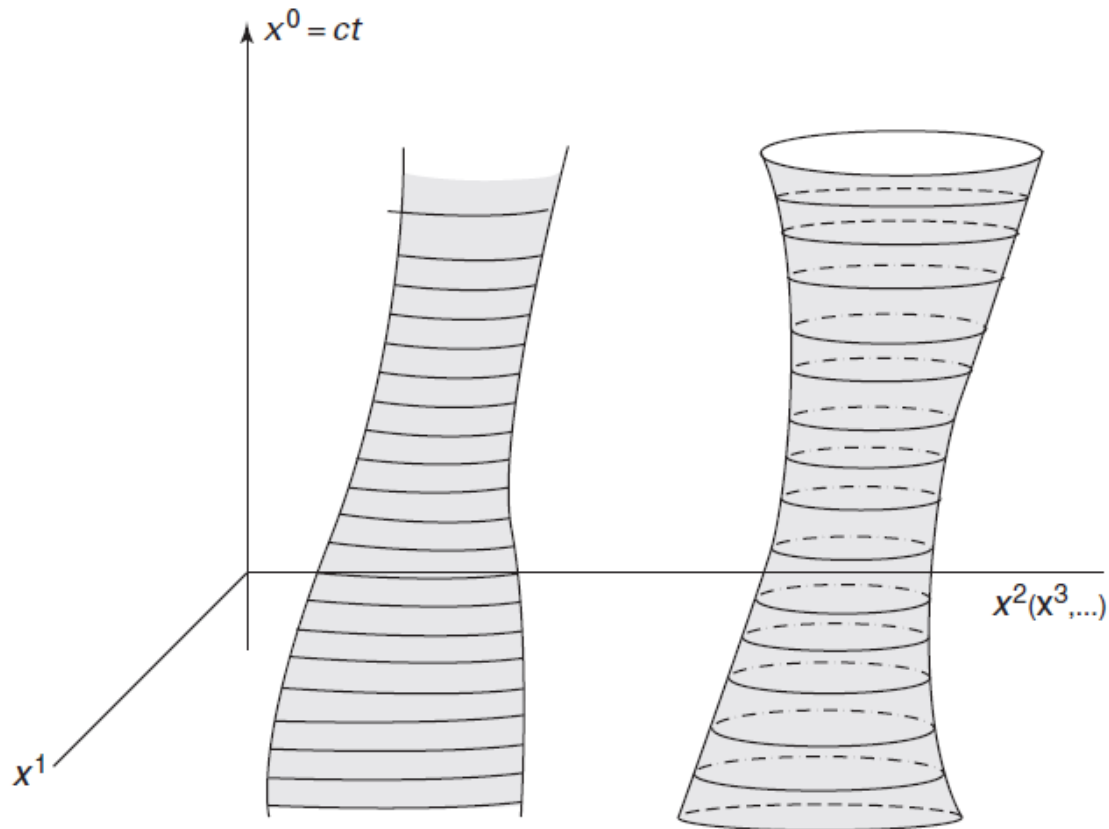
Relativistic Measure

Changing Variables

Lorentz Invariance

The relativistic action

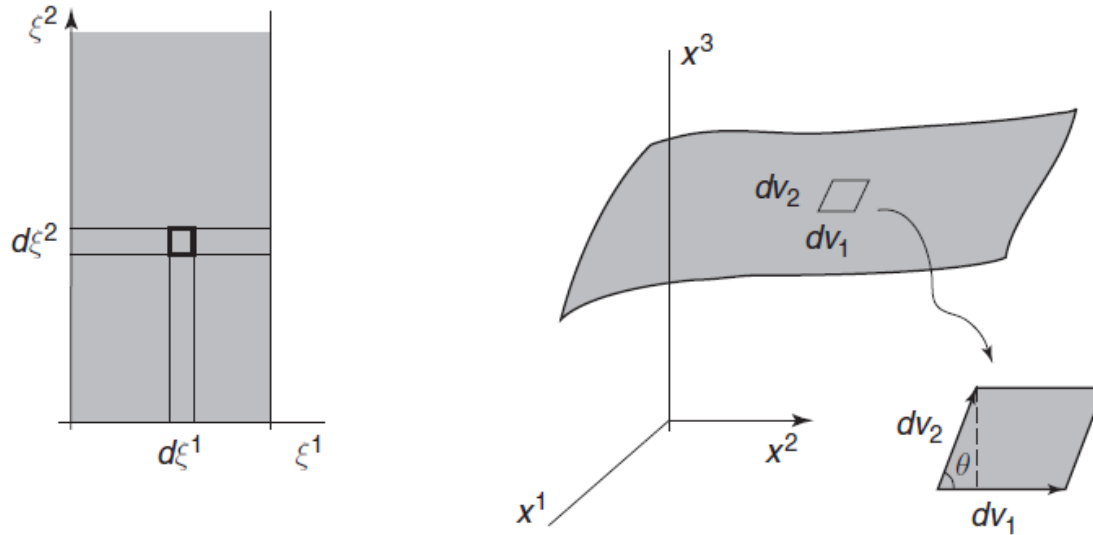
The relativistic string



ref: Zwiebach, p.101.

The relativistic action

The relativistic string



ref: Zwiebach, p.102.

$$\vec{x}(\xi^1, \xi^2) = \left(x^1(\xi^1, \xi^2), x^2(\xi^1, \xi^2), x^3(\xi^1, \xi^2) \right)$$

$$d\vec{v}_1 = \frac{\partial x}{\partial \xi^1} d\xi^1, \quad d\vec{v}_2 = \frac{\partial x}{\partial \xi^2} d\xi^2$$

The relativistic action

The relativistic string

$$\begin{aligned}dA &= |d\vec{v}_1| |d\vec{v}_2| |\sin \theta| = |d\vec{v}_1| |d\vec{v}_2| \sqrt{1 - \cos^2 \theta} \\ &= \sqrt{|d\vec{v}_1|^2 |d\vec{v}_2|^2 - |d\vec{v}_1|^2 |d\vec{v}_2|^2 \cos^2 \theta}\end{aligned}$$

$$dA = \sqrt{(d\vec{v}_1 \cdot d\vec{v}_1)(d\vec{v}_2 \cdot d\vec{v}_2) - (d\vec{v}_1 \cdot d\vec{v}_2)^2}$$

The relativistic action

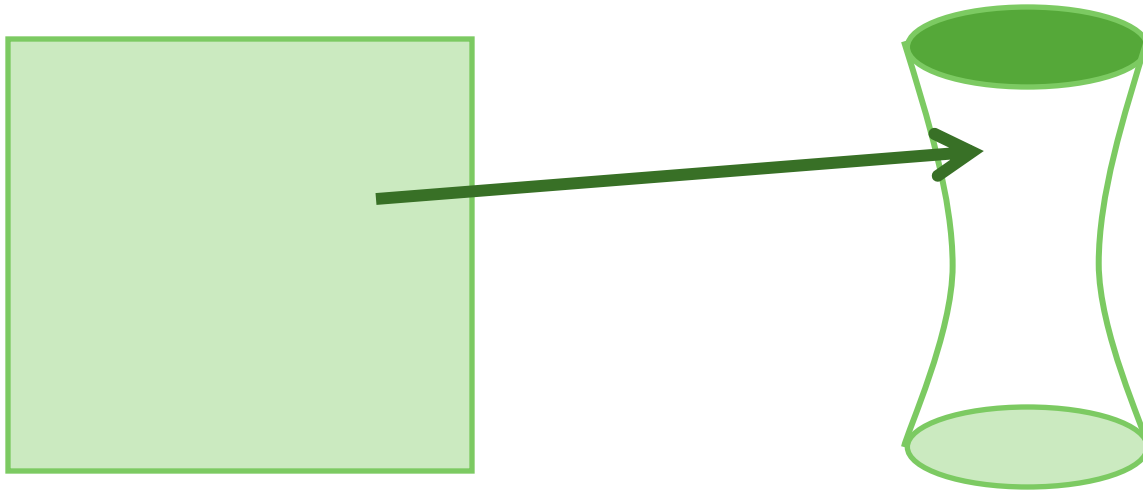
The relativistic string

$$dA = d\xi^1 d\xi^2 \sqrt{\left(\frac{\partial \vec{x}}{\partial \xi^1} \cdot \frac{\partial \vec{x}}{\partial \xi^1} \right) \left(\frac{\partial \vec{x}}{\partial \xi^2} \cdot \frac{\partial \vec{x}}{\partial \xi^2} \right) - \left(\frac{\partial \vec{x}}{\partial \xi^1} \cdot \frac{\partial \vec{x}}{\partial \xi^2} \right)^2}$$

$$A = \int d\xi^1 d\xi^2 \sqrt{\left(\frac{\partial \vec{x}}{\partial \xi^1} \cdot \frac{\partial \vec{x}}{\partial \xi^1} \right) \left(\frac{\partial \vec{x}}{\partial \xi^2} \cdot \frac{\partial \vec{x}}{\partial \xi^2} \right) - \left(\frac{\partial \vec{x}}{\partial \xi^1} \cdot \frac{\partial \vec{x}}{\partial \xi^2} \right)^2}$$

The relativistic action

The relativistic string



$$x^\mu = (x^0, x^1, \dots, x^d)$$

$$(\tau, \sigma) \rightarrow (X^0(\tau, \sigma), X^1(\tau, \sigma), \dots, X^d(\tau, \sigma))$$

The relativistic action

The relativistic string

$$A = \int d\tau d\sigma \sqrt{\left(\frac{\partial X}{\partial \tau} \cdot \frac{\partial X}{\partial \tau} \right) \left(\frac{\partial X}{\partial \sigma} \cdot \frac{\partial X}{\partial \sigma} \right) - \left(\frac{\partial X}{\partial \tau} \cdot \frac{\partial X}{\partial \sigma} \right)^2}$$

$$A = \int d\tau d\sigma \sqrt{\left(\frac{\partial X}{\partial \tau} \right)^2 \left(\frac{\partial X}{\partial \sigma} \right)^2 - \left(\frac{\partial X}{\partial \tau} \cdot \frac{\partial X}{\partial \sigma} \right)^2}$$

The relativistic action
The relativistic string

Worldline \rightarrow Worldsheet
 $dS \rightarrow dA$

$$S = -\frac{T_0}{c} \int_{\tau_i}^{\tau_f} d\tau \int_0^{\sigma_1} d\sigma \sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}$$

Relativistic String Action
Defined in Area on the Sheet

The relativistic action

The relativistic string

$$\mathcal{L}(\dot{X}^\mu, (X')^\mu) = -\frac{T_0}{c} \sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}$$

$$\mathcal{P}_\mu^\tau \equiv \frac{\partial \mathcal{L}}{\partial \dot{X}^\mu} = -\frac{T_0}{c} \frac{(\dot{X} \cdot X') X'_\mu - (X')^2 \dot{X}_\mu}{\sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}}$$

$$\mathcal{P}_\mu^\sigma \equiv \frac{\partial \mathcal{L}}{\partial X'^\mu} = -\frac{T_0}{c} \frac{(\dot{X} \cdot X') \dot{X}_\mu - (\dot{X})^2 X'_\mu}{\sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}}$$

The relativistic action

The relativistic string

$$\delta S = 0 \quad \rightarrow \quad \frac{\partial \mathcal{P}_\mu^\tau}{\partial \tau} + \frac{\partial \mathcal{P}_\mu^\sigma}{\partial \sigma} = 0$$

$$\frac{T_0}{c} = \frac{1}{2\pi\alpha'} \quad l_s = \sqrt{\alpha'}$$

The relativistic action

The relativistic string

$$\mathcal{L}(\dot{X}^\mu, (X')^\mu) = -\frac{T_0}{c} \sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}$$

$$\mathcal{P}_\mu^\tau \equiv \frac{\partial \mathcal{L}}{\partial \dot{X}^\mu} = -\frac{T_0}{c} \frac{(\dot{X} \cdot X') X'_\mu - (X')^2 \dot{X}_\mu}{\sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}}$$

$$\mathcal{P}_\mu^\sigma \equiv \frac{\partial \mathcal{L}}{\partial X'^\mu} = -\frac{T_0}{c} \frac{(\dot{X} \cdot X') \dot{X}_\mu - (\dot{X})^2 X'_\mu}{\sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}}$$

The relativistic action

The relativistic string

$$\delta S = 0 \quad \rightarrow \quad \frac{\partial \mathcal{P}_\mu^\tau}{\partial \tau} + \frac{\partial \mathcal{P}_\mu^\sigma}{\partial \sigma} = 0$$

$$\frac{T_0}{c} = \frac{1}{2\pi\alpha'} \quad l_s = \sqrt{\alpha'}$$

Effective parameterization: The light cone gauge

$$(n \cdot p) \sigma = \frac{2\pi}{\beta} \int_0^\sigma d\tilde{\sigma} n \cdot \mathcal{P}^\tau(\tau, \tilde{\sigma}) \quad : \sigma\text{-gauge}$$

$$n \cdot X(\tau, \sigma) = \beta \alpha' (n \cdot p) \tau \quad : \tau\text{-gauge}$$

$$\begin{aligned} \dot{X} \cdot X' &= 0 \\ \dot{X}^2 + X'^2 &= 0 \end{aligned}$$

Effective parameterization: The light cone gauge

A bunch of fancy $\frac{1}{2\pi\alpha'}$ math stuff...

The Wave Equation... $\frac{\partial \mathcal{P}^\tau}{\partial \tau}$ $\frac{\partial \mathcal{P}^\sigma}{\partial \sigma}$ AGAIN!!

Solving the wave equation

Where is the string going?

$$X^\mu(\tau, \sigma) = x_0^\mu + \sqrt{2\alpha'}\alpha_0^\mu\tau - i\sqrt{2\alpha'}\sum_{n\neq 0}\frac{1}{n}\alpha_n^\mu e^{-in\tau}\cos n\sigma$$

α_0^μ How is it oscillating? n

Becoming particles

The string energy stored

$$M^2 = \frac{1}{\alpha'} \sum_{n=1}^{\infty} n a_n^{I*} a_n^I$$

translates into particle mass

Closed strings and open strings

Closed strings → Gravitons

Open strings → Other bosons

Acknowledgements

Thank You:

Dr. Melanie Becker

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Sean Downes



Acknowledgements

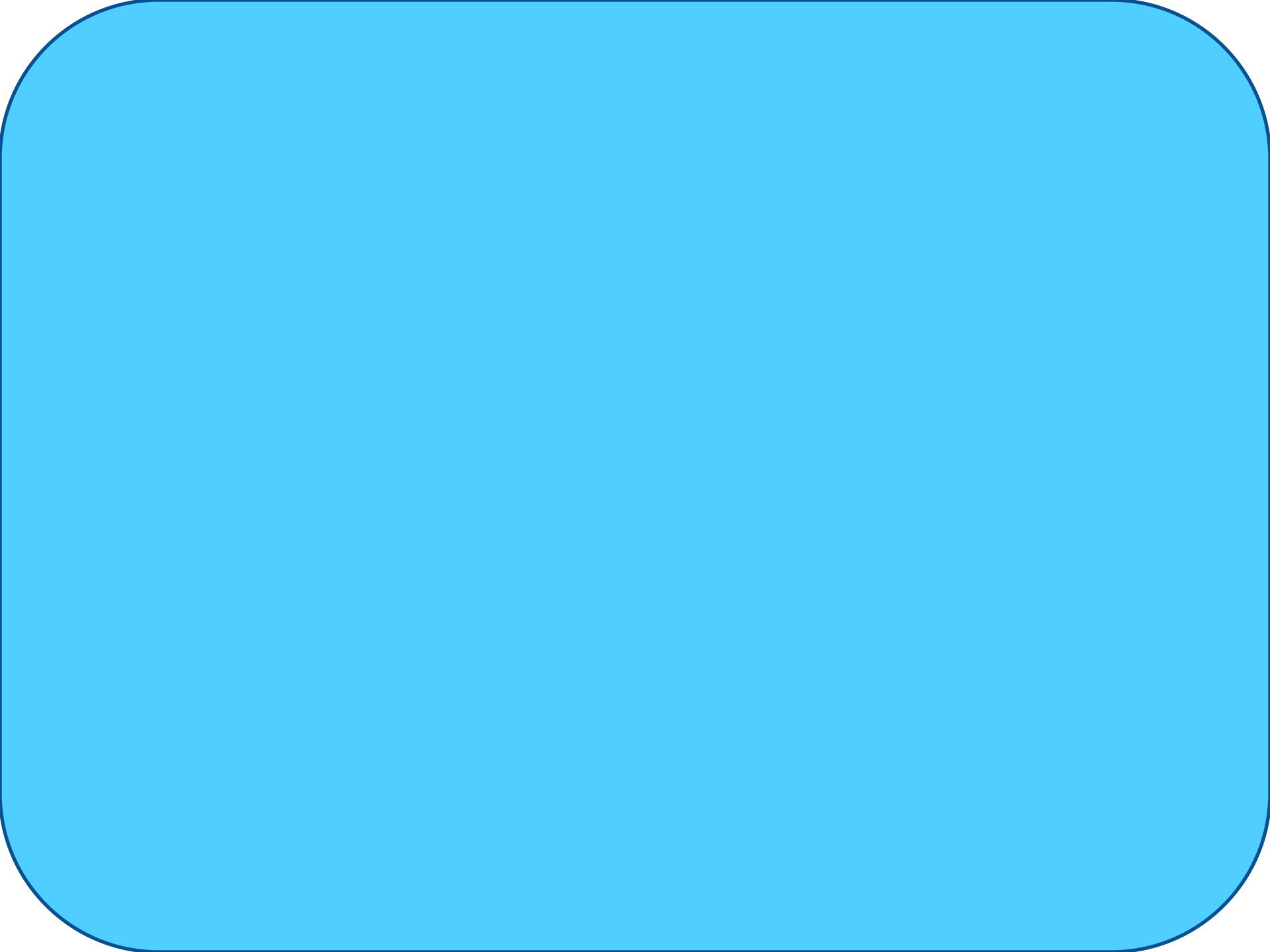


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The relativistic action

The relativistic point particle

$$\delta S = - \int_{\tau_i}^{\tau_f} d\tau \delta x^\mu(\tau) \frac{dp_\mu}{d\tau}$$

$$\delta S = 0 \quad \rightarrow \quad \frac{dp_\mu}{d\tau} = 0$$

$$\frac{dp_\mu}{d\tau} = \frac{d}{d\tau}(p_\mu) = \frac{d}{d\tau}(\eta_{\mu\nu} p^\nu) = \eta_{\mu\nu} \frac{d}{d\tau}(p^\nu)$$

$$p^\mu = mc \frac{dx^\mu}{ds} \quad \rightarrow \quad \frac{d^2 x^\mu}{ds^2} = 0$$